

Coupled Oscillations and Normal Modes.

$$1^{\circ} \quad L = T - V$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} = \frac{\partial L}{\partial q_i} \quad (i = 1, 2, \dots, N) \quad n \text{ particles.}$$

2^o ~~guess~~ guess normal mode exists, substitute x_i with $A_i \cos(\omega t - \delta)$
then get a ^{set of} linear equations of A_i .

3^o ~~let~~ let the coefficient ~~determinant~~ determinant be zero to get the non-trivial solution according to different ω .

4^o the general solution will be the linear combination of the n normal modes.

General Theory of Vibrating Systems.

1^o Now we discuss the general case that a system with n degrees of freedom that is oscillating about a point of equilibrium.

We specify the configuration by n generalized coordinates q_1, q_2, \dots, q_n .

$$\vec{q} = (q_1, q_2, \dots, q_n)$$

we assume the system is conservative, and it has a potential energy.

$$U(q_1, q_2, \dots, q_n) = U(\vec{q})$$

we assume the constraint here is constant, hence the displacement of the α -th particle is. $\vec{r}_\alpha = \vec{r}_\alpha(q_1, q_2, \dots, q_n)$

Then T can be written in the form:

$$T = T(\vec{q}, \dot{\vec{q}}) = \frac{1}{2} \sum_{j,k} A_{jk}(\vec{q}) \dot{q}_j \dot{q}_k$$

2^o The Taylor Expansion of $U(\vec{q})$ is
(approximation)

$$U(\vec{q}) = U(\omega) + \sum_j \frac{\partial U}{\partial q_j} q_j + \frac{1}{2} \sum_{j,k} \frac{\partial^2 U}{\partial q_j \partial q_k} q_j q_k + \dots$$

let $U(\omega) = 0$, and for $\vec{q} = 0$ is the equilibrium point, hence

$$\frac{\partial U}{\partial q_j}(\omega) = 0 \text{ for any } j.$$

We ~~now~~ now take the second order terms only, and get

$$U(\vec{q}) = \frac{1}{2} \sum_{j,k} K_{jk} q_j q_k.$$

let $\vec{q} = 0$, and the T is just the ~~second~~ second order terms:

$$T(\vec{q}) = \frac{1}{2} \sum_{j,k} M_{jk} \dot{q}_j \dot{q}_k, \text{ where } M_{jk} = A_{jk}(0).$$

Hence the Lagrangian,

$$\mathcal{L}(\vec{q}, \dot{\vec{q}}) = T(\vec{q}) - U(\vec{q}).$$

$$3^{\circ} \quad \frac{\partial V}{\partial q_i} = \sum_j k_{ij} q_j \quad (i=1, \dots, n).$$

(Equation of Motion in Matrix Form).

$$\frac{d}{dt} \frac{\partial T}{\partial \dot{q}_i} = \sum_j M_{ij} \ddot{q}_j \quad (i=1, \dots, n).$$

Hence, the Lagrange's Equations are

$$\sum_j M_{ij} \ddot{q}_j = - \sum_j k_{ij} q_j \quad [i=1, 2, \dots, n]$$

in Matrix form is, $M\ddot{\vec{q}} = -K\vec{q}$, where M and K are "mass" and "spring - constraint" matrices.

$$4^{\circ} \quad \text{try solutions in the form of } \vec{q} = \vec{a} \cos(\omega t - \delta),$$

(Method of Normal mode) we get the eigenvalue equation,

$$(K - \omega^2 M) \vec{a} = 0$$

then let ω satisfies the characteristic (or ~~sec~~ secular) equation,

$$\det(K - \omega^2 M) = 0$$

if ~~then~~ we get n different ~~non~~ non-negative ω , and \vec{a} respectively, the general solution is just the linear combination of the normal mode solutions.